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Porphyrin Core Expansion and Doming in Heme Proteins. New Evidence from Resonance Raman Spectra of Six-Coordinate High-Spin Iron(III) Hemes

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Abstract: Resonance Raman spectra are reported for high-spin bis(dimethyl sulfoxide) (Me₂SO)₂Fe¹¹¹ complexes of protoporphyrin 1X dimethyl ester (PP), octaethylporphyrin (OEP), and tetraphenylporphyrin (TPP), and for aquo- and fluoromethemoglobin (Hb) and myoglobin (Mb), with emphasis on the three bands (II, IV, and V) which have previously been shown to be sensitive to iron spin state. These frequencies are essentially the same for [(Me₂SO)₂Fe^{III}PP]⁺ as for aquo-metHb and Mb, which had previously been thought to be anomalously low. Since recent crystal structure determinations have shown the highspin bisaquo and bis(tetramethylene sulfoxide) complexes of Fe¹¹¹TPP to be planar porphyrins, with expanded cores, this experiment strongly supports core expansion as the determinant of the metHb and Mb spin-marker frequencies, and renders the hypothesis of protein-induced doming unlikely. Since protein crystallography has placed the iron atom 0.40 Å from the mean heme plane in metMb and 0.23 and 0.07 Å from the mean heme plane in the β and α chains of metHb, the lack of any associated variation in the spin-marker frequencies indicates that the core size is not determined by the disposition of the iron atom, but probably by steric interactions of the axial ligands with the pyrrole nitrogen atoms. Binding of fluoride to metHb and Mb leaves bands II and V unshifted, but lowers band IV by 5 cm⁻¹. The similarity in spin-marker frequencies between native horseradish peroxidase (HRP) or cytochrome c' and the previously studied five-coordinate Fe¹¹¹ hemes indicates that these proteins contain five- rather than six-coordinate Fe¹¹¹ hemes; in the case of HRP six coordination is accessible with exogenous ligands. It is suggested that the intermediate-spin state observed for these proteins results from weakening of the bond between Fe¹¹¹ and the single axial ligand. The correlation with core size of the spin-marker frequencies is reexamined. Nonplanar hemes are found to have frequencies that are significantly depressed, presumably due to loss of π conjugation at the methine bridge. The relation between pyrrole tilt and the methine dihedral angles is derived for both ruffled and domed hemes, and is used to develop a simple relation that satisfactorily reproduces the spin-marker frequencies for both planar and nonplanar hemes. This model gives reasonable estimates of the σ and π contributions to the methine bond-stretching force constant. It is concluded that both core expansion and pyrrole tilting contribute to frequency lowerings, with core expansion dominant for moderate tilt angles. There is no evidence in any of the heme proteins so far studied of extra doming induced by the protein.

Introduction

Resonance Raman spectra of heme proteins contain a rich assortment of porphyrin vibrational modes, whose frequencies have been well catalogued.¹⁻⁷ Some of the modes are sensitive to chemical alterations at the central iron atom, and are therefore of great interest as potential structure monitors. In an early study² it was observed that one set of bands was sensitive to changes in oxidation state of the iron atom while another set was sensitive to changes in its spin state. The oxidation state marker bands were suggested^{2,7a} to be responding to changes in porphyrin π^* orbital occupancy via back-donation from the iron d_{π} orbitals. This interpretation was subsequently confirmed⁸ with a graded series of π -acceptor axial ligands, which compete with the porphyrin ring for the iron d_{π} electrons. An interesting consequence of this mechanism is that π -donor ligands can shift the oxidation state marker

frequencies to values below the typical ones associated with nonacceptor ligands, as has recently been observed by Champion and Gunsalus⁶ and by Kitagawa et al.^{7b} for cytochrome P₄₅₀, which is believed to have cysteine mercaptide, a π donor, as an axial ligand.

With respect to the spin-state marker bands our interpretation² was that the frequency decreases observed for high-spin hemes resulted from a doming of the porphyrin ring accompanying the well-known out-of-plane displacement of the iron atom toward the proximal ligand.9 The pyrrole rings are expected to tilt, in order to maintain overlap with the iron orbitals, but this would be at the expense of π conjugation at the porphyrin methine bridges, resulting in the observed frequency lowerings. A subsequent normal-coordinate calculation^{10a} indicated that this mechanism could account for the experimental data. A discrepancy between the spin marker frequencies of aquomethemoglobin (metHb) and of fivecoordinate high-spin Fe¹¹¹ heme suggested an enhanced degree of doming for this form of the protein.⁸

A different interpretation of the spin marker band (band IV) at ~1580 cm⁻¹, which is easily recognized via its anomalous polarization,^{10b} was advanced by Spaulding et al.,¹¹ who observed large decreases in this frequency for planar porphyrins with large atoms at their center, such as Sn^{1V} and Ag¹¹. To accommodate these atoms the porphyrin core is forced to expand, and Spaulding et al. found an inverse correlation between the band IV frequency and Ct-N, the distance from the center of the porphyrin to the pyrrole nitrogen atoms. The effect derives from stretching and bending of the methine bridge bonds which is observed to correlate with Ct-N,^{9,11} Warshel's potential energy calculations¹² confirm that porphyrin core expansion is accommodated by methine bridge bond stretching and deformation.

Spaulding et al.¹¹ suggested that core expansion was also responsible for the band IV frequency shifts in heme proteins, and that the out-of-plane displacement of the iron atom could be calculated by triangulation, on the assumption of fixed radii for high- and low-spin iron(11) and -(111). They concluded that the anomaly between metHb and five-coordinate Fe¹¹¹ porphyrin required an extra expansion of the porphyrin core in the protein, possibly induced by coordination of the sixth ligand, H_2O , which might pull the iron atom closer to the heme plane. At the time we believed this inference to be implausible, since we expected⁸ that movement of the iron atom into the porphyrin ring would be accommodated by emptying of the inplane $d_{x^2-y^2}$ orbital, via spin pairing, and a resultant contraction, rather than expansion of the porphyrin core. This effect had been established for iron(11) tetraphenylporphine (TPP), which, in the absence of any axial ligand, adopts a structure with the iron atom in the center of a contracted porphyrin core; the depopulation of $d_{x^2-y^2}$ is reflected in an intermediate spin state.13

Subsequent events, however, have shown this reasoning to be inapplicable to high-spin Fe¹¹¹ with two axial ligands. Refinement of the metHb crystal structure¹⁴ has revised downward the estimate of the out-of-plane iron displacement to 0.23 and 0.07 Å for the β and α chains. And recently the bis(dimethyl sulfoxide) adduct of Fe¹¹¹TPP has been shown by NMR to be high spin in solution,¹⁵ while the crystal structure of the high-spin bisaquo¹⁶ and bis(tetramethylene sulfoxide)¹⁷ complexes of Fe¹¹¹TPP reveal an in-plane iron atom with an expanded porphyrin core. These crystallographic findings provide an important test of the core expansion hypothesis, which requires that planar six-coordinate high-spin Fe¹¹¹ hemes show spin marker frequencies appreciably lower than outof-plane five-coordinate derivatives. The results of the present study confirm this prediction. They establish that six coordination is a sufficient explanation of the low spin-marker frequencies observed for metHb. At the same time, the Raman data provide convincing evidence that native horseradish peroxidase (HRP) and cytochrome c' are five coordinate.

In the light of these results and of the available data on nonplanar porphyrins, we have reexamined the doming hypothesis. We find that nonplanarity induces a calculable additional lowering of the spin-marker frequencies. While not as influential as core size, the effect is not negligible. A given set of spin-marker frequencies is consistent with a range of values for the core size and the pyrrole tilt angle.

Experimental Section

Raman spectra were obtained using radiation from either a Spectra Physics 170 Ar^+ laser or a Coherent Radiation 590 dye laser, and a Spex 1401 double monochromator equipped with Ortec/digital photon counting electronics. The spectrometer and photon counting were under control of a Data General Nova 3 computer. Electron para-



Figure 1. Electronic absorption spectra of (A) $[(Me_2SO)_2Fe^{11}PP]^+$, (B) aquometHb, (C) ClO₄Fe¹¹PP. All concentrations are 0.1 mM, except for Soret spectra, which were 0.01 mM.

magnetic resonance spectra were recorded on a Varian E-12 spectrometer. Electronic absorption spectra were recorded with a Cary 118 ratio recording spectrophotometer.

Hemin chloride (Sigma) was converted to the dimethyl ester¹⁸ and then to the μ -oxo dimer by the method of Sadasivan et al.¹⁹ Hemin perchlorate was prepared by shaking the oxo dimer in methylene chloride with two volumes of 4 M HClO₄. The organic phase was separated and anhydrous Na2SO4 added. The mixture was swirled vigorously for 15 min and filtered. The filtrate was then allowed to stand over anhydrous Na_2SO_4 for 5 days. It is extremely important that all water be removed. Traces of water hasten decomposition and interfere with the subsequent formation of the dimethyl sulfoxide (Me₂SO) complex. lron(III) octaethylporphyrin perchlorate was prepared as described by Dolphin et al.²² Methylene chloride was washed first with saturated Na2CO3, then twice with water, and dried over Na₂SO₄. Spectral grade Me₂SO (Baker) was purified by standing over 4 Å molecular sieves. Hemoglobin was obtained from lysed human erythrocytes. It was converted to the Fe¹¹¹ form by gentle stirring with a threefold molar excess of potassium ferricyanide, followed by chromatography on a 3×20 cm column of Sephadex G-25 coarse, equilibrated with 0.05 M phosphate, pH 6.6 Horse myoglobin was obtained from Miles Research and purified by the method of Hapner et al.20

Titrations were performed by making additions of titrant (Me₂SO) directly to a cuvette containing hemin perchlorate in methylene chloride. The absorption spectrum was recorded and the resulting changes in absorption were used to determine the stoichiometry of the reaction by well-established methods.²¹ The concentration of the hemin perchlorate solution was determined by titration with imidazole, assuming the formation of the bisimidazole complex to be quantitative.

Results and Discussion

Characterization of Bis-Me₂SO Fe^{III} Hemes. A methylene chloride solution of (ClO₄⁻)Fe^{III}PP (PP = protoporphyrin IX dimethyl ester) had a visible absorption spectrum consistent with that reported for the OEP (octaethylporphyrin) analogue.²² Addition of dimethyl sulfoxide (Me₂SO) shifted the absorption to a new spectrum as shown in Figure 1. Spectrophotometric titration establishes a 2;1 stoichiometry for the new complex, with quantitative addition of the first Me₂SO molecule. A second Me₂SO molecule is added, with $K = 23.5 \pm 3.4 \text{ M}^{-1}$. ESR spectroscopy of the solution with excess Me₂SO gave a typical high-spin Fe^{III} spectrum with $g \approx 6$. The analogous bis-Me₂SO Fe^{III}TPP complex has been characterized by La Mar and co-workers,¹⁵ via NMR titration, which



Figure 2. Resonance Raman spectra of Fe^{III} heme derivatives and proteins obtained with 514.5-nm excitation (except for $[(Me_2SO)_2Fe^{III}PP]^+$, which was excited at 5017 Å). All samples were ~1 mM. Conditions: laser power 300 mW, slit width 3 cm⁻¹, scanning speed 3 s/step (0.5 cm⁻¹/step). Spectra shown are the accumulation of six complete scans and subjected to a nine-point quintic-quartic smooth. Frequencies are relative to the 1423-cm⁻¹ line of methylene chloride. The dotted lines mark the band IV and V positions for fluorometHb.

also established the equivalence of the two bound Me₂SO ligands in the slow exchange temperature range, consistent with the iron atom being in the porphyrin plane. Our solution of $(ClO_4^-)Fe^{111}$ TPP with excess Me₂SO also displayed a highspin Fe¹¹¹ EPR spectrum. A planar structure has been established by X-ray crystallography for the bisaquo¹⁶ and bis(tetramethylene sulfoxide)¹⁷ complexes of Fe¹¹¹TPP. It is most likely that high-spin Fe¹¹¹ bis-Me₂SO complexes are also planar. for porphyrins other than TPP, although this has yet to be confirmed by X-ray crystallography.

Spin-Marker Raman Frequencies. Figure 2 shows RR spectra in the 1500-1650-cm⁻¹ region for [(Me₂SO)₂-Fe^{III}PP]⁺, (Cl⁻)Fe^{III}PP, and aquo- and fluorometHb (hemoglobin) and Mb (myoglobin). The observed frequencies and polarizations are given in Table 1. Table 11 lists the frequencies of the three bands, labeled 11, 1V, and V, which have been found to be sensitive to variations in the spin state of the iron atom.^{2,8} Included in this table are representative examples of other spin and oxidation states, and a series of high-spin heme protein derivatives. The last column contains values of C_t -N, the porphyrin center to pyrrole nitrogen distance, which is a measure of the porphyrin core size. Some of the listed distances are obtained from crystal structures for porphyrins, including TPP, other than those for which the Raman frequencies are given. Structural parameters do not seem to vary significantly, however, with alterations in the porphyrin substituents.⁹ For comparison of vibrational frequencies it is essential that the methine bridges have hydrogen atoms and that the β -pyrrole positions have carbon substituents.^{2,8}

Table 11 makes it evident that the general correlation of the band IV frequency with core size observed by Spaulding et al.¹¹ holds for all the iron hemes, and it also holds for bands 11 and V (although the low-spin Fe^{II} frequencies are exceptional—see below). For comparable spin states the Fe-pyrrole bonds are slightly longer for Fe¹¹ than Fe¹¹¹, reflecting the reduced effective charge, although the shortest bonds are found for four-coordinated (intermediate-spin) Fe^{II}, the effective charge being increased by the absence of any axial ligands.¹³ Highspin hemes have longer Fe-pyrrole bonds than low-spin hemes because of the partial occupancy of the antibonding $d_{x^2-v^2}$ orbitals. For five-coordinate high-spin hemes this repulsion is partially relieved by the iron atom moving out of the porphyrin plane in the direction of the axial ligand. There is also a significant expansion of the porphyrin core, however. When high-spin Fe¹¹¹ is forced into the plane by two equivalent axial ligands the porphyrin core expands further.

 $[(Me_2SO)_2Fe^{111}PP]^+$ gives RR frequencies (Table 1) which are close to those found in other iron TPP complexes.²³ Vibrational coupling with ring substituents produces a substantially altered normal mode pattern for TPPs, and clear-cut spin- and oxidation-state markers are not observed.²³ There are, however, two polarized bands, at ~1550 and ~1360 cm⁻¹ (also one at ~390 cm⁻¹, in a frequency region not included in the present study) which vary systematically with both oxidation and spin state.²³ They are found at 1568 and 1370 cm⁻¹ for low-spin Fe¹¹¹, and at 1555 and 1366 cm⁻¹ for five-coordinate high-spin Fe¹¹¹. The present data place the six-coordinate high-spin Fe¹¹¹ frequencies at 1550 and 1360 cm⁻¹, below those of five-coordinate Fe¹¹¹, as expected from core expansion.

Insensitivity to Fe Out-of-Plane Displacement. MetHb and Mb. Spaulding et al.¹¹ proposed that, if C_t-N could be determined from the band IV frequency, then the distance of the iron atom from the plane of the four nitrogen atoms could be calculated by triangulation, assuming that the Fe-pyrrole distance for high-spin Fe^{III} or Fe^{III} is constant. The crystal structure of the planar high-spin Fe^{III}TPP derivatives^{16,17} shows this assumption to be incorrect. Their Fe-pyrrole distances are shorter than those found for five-coordinate high-spin Fe^{III} complexes.⁹ Scholler and Hoffman²⁴ have proposed a different method for estimating the frequency shift accompanying out-of-plane displacement.

The present RR data on metHb and Mb, taken together with the X-ray crystallographic results,^{14,25} establish that *core expansion and iron out-of-plane displacement are unrelated*, at least for six-coordinate high-spin Fe^{III}. The frequencies of aquometHb and Mb are the same despite substantially dif-

Table I. Raman Frequencies (cm⁻¹) of Six-Coordinate Ferric Hemes

aquometHb	1481 (p)	1549 <i>ª</i>	1561 (ap)	1583 (p)	1610 (dp)	1623 <i>ª</i>
aquometMb	1482 (p)	1546	1562 (ap)	1582 (p)	1611 (dp)	1626
fluorometHb	1482 (p)	1546	1556 (ap)	1583 (p)	1608 (dp)	1621
fluorometMb	1482 (p)	1544	1557 (ap)	1583 (p)	1609 (dp)	1619
[(Me ₂ SO) ₂ Fe ¹¹¹ PP] ⁺	1475 (p)	1545	1560 (ap)	1583 (p)	1610 (dp)	1628
[(Me ₂ SO) ₂ Fe ^{III} OEP] ⁺	1481 (p)	1552	1563 (ap)	1577 (p)	1613 (dp)	1629
$[(Me_2SO)_2Fe^{III}TPP]^+$	1333 (ap)	1360 (p)	1491 (p)	1508 (ap)	1550 (p)	1597 (p)

^a These bands appear as shoulders, and are of uncertain polarization.

Table II. Raman Frequencies (cm⁻¹) of Structure-Sensitive Bands and C_t -N Distances (Å)

	11	ıv	V	C ₁ -N	ref
Fe ¹¹ intermediate spin Fe ¹¹ MP	1506	1589	1642	1.97213	8
Fe ¹¹¹ low-spin [(1m) ₂ Fe ¹¹¹ MP] ⁺	1505	1584	1640	1.989 ^{54a}	8.11
Fe ¹¹ low-spin [(1m) ₂ Fe ¹¹ PP]	1491	1583	1617	2.004 ^{54b}	8,11
Fe ¹¹¹ high-spin 5c ClFe ¹¹¹ PP	1495	1572	1632	2.0199	11
Fe ^{III} high-spin 6c [(Me ₂ SO) ₂ Fe ^{III} PP] ⁺	1475	1560	1610	2.04517	this work, 11
Fe ^{III} high-spin 6c (Me ₂ SO) ₂ Fe ^{III} OEP ⁺	1481	1563	1613		this work
Fe ¹¹ high-spin 2-MelmFe ¹¹ MP	1472	1558	1606	2.04456	8.11
high-spin Fe ¹¹ proteins:					
dcoxyHb	1473	1552	1607		
Fe ¹¹ HRP	1472	1553	1605		
high-spin Fe ¹¹¹ proteins					
aquometHb	1481	1561	1610		this work
aquometMb	1482	1562	1611		this work
fluorometHb	1482	1556	1608		this work
fluoromet Mb	1482	1557	1609		this work
fluoro-HRP	1482	1555	1608		32
native HRP	1500	1575	1608		32
eyiochrome c' (pH 6.9)	1500	1578	1637		42
cytochrome c' (pH 10.3)	1496	1573	1633		42

ferent Fe-mean heme plane displacements: 0.40 Å for metMb²⁵ and 0.23 and 0.07 Å for the β and α chains of metHb.¹⁴ Moreover, the metHb RR spectrum (Figure 2) shows no detectable splitting of the spin-marker band reflecting the chain inequivalence. The frequencies for both proteins are quite close to those of $[(Me_2SO)_2Fe^{111}PP]^+$ (Table 11), but substantially lower than those of $(Cl^-)Fe^{111}PP$ or other fivecoordinate high-spin hemes. It appears that six coordination per se is the main determinant of the spin-marker frequencies, and, by implication, of the core expansion.

This conclusion implies that the porphyrin core size is not determined by the size of the high-spin ferric ion itself, but rather by interactions with the axial ligands. Recent ab initio calculations by Olafson and Goddard²⁶ and semiempirical calculations by Warshel²⁷ point to the importance of nonbonded repulsions between the axial ligands and the pyrrole nitrogen atoms in the structure and energetics of the heme group. For five-coordinate hemes, it appears²⁶ that the outof-plane displacement of the iron atom depends on the optimization of iron-ligand bonding within a coordination geometry that is fixed by the axial ligand-pyrrole interactions. This ordering of energies is likely to hold also for six-coordinate high-spin hemes. When the axial ligands are equivalent, then the iron atom is constrained to lie in the heme plane. The potential curve for out-of-plane displacement is likely to be relatively flat, however.²⁷ When the ligands are not equivalent, the position of the iron atom can be expected to depend on the relative bond strengths to the axial ligands and their nonbonded interactions with the pyrrole rings. The balance of these forces may provide for considerable flexibility in the axial geometry. It is of interest in this connection that in (py)(Cl⁻)Mn¹¹¹TPP²⁸ the Mn atom is displaced toward the chloride ion and the long bond is to pyridine. Although pyridine is a stronger field ligand, this is outweighed for the high-spin Mn^{3+} ion by the negative charge on Cl⁻.

metHb and metMb have the same axial ligand set, imidazole and water. For the metHb α chains the axial forces are evenly balanced and the iron atom is only 0.07 Å from the mean heme plane. For the β chains and for metMb tertiary forces apparently pull on the imidazole, the out-of-plane displacement increasing to 0.23 and 0.42 Å respectively. Presumably the imidazole-pyrrole repulsion is weakened, but the water-pyrrole repulsion must be strengthened by the same amount since the RR frequencies show the porphyrin core size to remain constant. The overall change in the energy of the heme complex may be quite small. It must be sufficient, however, to account for the lower fraction of low-spin forms in metMb complexes relative to the corresponding metHb complexes.²⁹ As Ladner et al.¹⁴ have pointed out, this trend is consistent with the larger out-of-plane displacement observed for metMb. The potential curve for out-of-plane displacement is expected to be steeper for low-spin than for high-spin Fe¹¹¹ because of the greater bond directionality of the former. Consequently, protein forces favoring out-of-plane displacement will raise the energy of the low-spin form relative to the high-spin form.

The fluoride complexes of metHb and Mb (Figure 2 and Table 11) show essentially the same band 11 and V frequencies as do the aquo complexes, but band 1V is 5 cm⁻¹ lower. A difference Fourier map of fluoro- vs. aquo metHb, at 2.8-Å resolution, shows some tertiary structure alterations but no significant changes in the heme structure.³⁰ The decreased band 1V frequency may reflect these tertiary alterations, or it may result from increased nonbonded interactions of the anionic fluoride ligands.

Five Coordination in Horseradish Peroxidase and Cytochrome c'. As noted previously,^{8,31} the spin-marker frequencies for native horseradish peroxidase (HRP, which contains high-spin Fe¹¹¹), 1500, 1575, and 1630 cm⁻¹ for bands 11, IV, and V, are close to those of five-coordinate Fe¹¹¹ porphyrins, rather than those of metHb or metMb (see Table 11). In view of the present evidence that the metHb and metMb frequencies are characteristic of six-coordinate high-spin Fe¹¹¹ porphyrins, the HRP frequencies clearly point to five coordination for the native form of the protein. This inference is consistent with the



Figure 3. Frequency vs. C_t -N plots for bands 11, 1V, and V. The filled circles, which are used to establish the straight lines, are data taken from ref 11. Open circles are data points for the indicated nonplanar porphyrins:^{8,11} (2-Melm)Fe^{II}MP, (Cl⁻)Fe^{III}MP, [(1n)₂Fe^{III}MP]⁺, Fe^{II}MP, and Ni^{II}(OEP) (D_{2d}), with structural data taken from ref 56, 59, 54, 13, and 53, respectively. Crosses are frequencies calculated via eq 2. Open squares are planar porphyrins, (Im)₂Fe^{II}(MP) and [(Me₂SO)Fe^{III}PP]⁺, which give discrepant band 11 and V frequencies, as discussed in the text.

NMR relaxation data presented by Lanir and Shejter,³² which establish that the native HRP heme group does *not* have an exchangeable water molecule bound to it. A five-coordinate heme would have no bound water if the fifth ligand were provided by the protein. The identity of the fifth ligand is uncertain, but imidazole is implicated by ultraviolet difference spectra³³ and by the ESR hyperfine splitting pattern in the nitrosyl complex of reduced HRP.³⁴

While solvent water does not bind to HRP, exogenous anionic ligands can be made to bind. The fluoride complex gives RR frequencies³¹ which are the same as those of fluorometHb (Table 11) implying six coordination, while the CN⁻ complex is low spin and six coordinate.³¹ However, anion binding is always accompanied by the binding of a proton to the protein.35 This has been interpreted by George and co-workers³⁶ as requiring displacement of a deprotonated ligand, such as tyrosine phenolate, from the heme iron atom. The present data, pointing to five coordination, suggest that the deprotonated group is not actually bound to the iron atom, but is close enough to block access to it. Alternatively a deprotonated distal group is bound to the heme iron, but the bond to the proximal imidazole is broken, being formed upon binding of exogenous ligands which displace the distal group. (The spectroscopic evidence bearing on imidazole coordination does not establish an intact bond in native HRP). In view of the flexibility of high-spin Fe¹¹¹ with respect to the porphyrin core (inversion of five-coordinate Fe¹¹¹ through the porphyrin is known to occur on a time scale of $\sim 10^{-2} \text{ s}^{37}$) switching of endogenous ligands is not implausible. Just this structure has been found by Pulsinelli et al.³⁸ in the Hb M Boston valency hybrid, the Fe³⁺ ion of the mutant α chains being bound to the distal tyrosine and not to the proximal imidazole. It would be of interest to examine the RR spectrum of this hybrid, since we expect the contribution from the Fe¹¹¹ hemes to resemble native HRP rather than metHb.

Spaulding et al.¹¹ reported that the benzhydroxamate complex of HRP gives a band IV frequency of 1566 cm⁻¹,

intermediate between the frequencies observed for the fluoride (1555 cm⁻¹) and for native HRP (1575 cm⁻¹). It may be that binding of benzhydroxamate induces a weak axial interaction on the distal side, either with itself or with solvent water. The latter possibility is suggested by the ESR study of the hydroxamate complex by Gupta and Schonbaum,³⁹ which established a definite hyperfine splitting by ¹⁷OH₂ solvent, although the extent of the splitting (1 G) was much smaller than that observed for metMb (18 G).⁴⁰

Another Fe¹¹¹ heme protein that shows a five-coordinate RR pattern is cytochrome c' from *Rhodopseudomonas palust*ris.^{41,42} In this case the heme iron atom is not accessible to exogenous ligands in the oxidized form of the protein.⁴³ The magnetic moment is intermediate between low- and high-spin Fe¹¹¹,⁴⁴ but the Raman data rule out a spin-state mixture.^{41,42} EPR spectroscopy has confirmed that the spin state itself is intermediate.⁴⁵ Also HRP, although high spin, is believed to be poised at the edge of an intermediate spin state, which is stabilized at low temperature.⁴⁶

In this connection it is of considerable interest that the perchlorate complexes of Fe¹¹¹OEP²² and Fe¹¹¹TPP¹⁶ have been shown to be intermediate spin. The crystal structure of the TPP complex¹⁶ shows it to be five coordinate, but the Fepyrrole bonds are nearly as short as those found in low-spin Fe^{III} porphyrin consistent with depopulation of the antibonding $d_{x^2-y^2}$ orbital.¹⁶ Perchlorate is a very weak ligand field, and thus it appears that, if the axial field is weakened sufficiently, then the intermediate-spin state is stabilized. This suggests that a protein containing a five-coordinate Fe¹¹¹ heme can convert from high to intermediate spin by stretching the Fe axial ligand bond. The EPR work of Maltempo et al.⁴⁵ suggests that the transition is continuous, rather than discrete, since they find a quantum admixture of high- and intermediate-spin contributions, i.e., a state "intermediate" between $S = \frac{5}{2}$ and S = $^{3}/_{2}$.

The structural parameters for $(ClO_4^-)Fe^{111}TPP^{16}$ indicate that stabilization of the intermediate-spin state should be accompanied by a contraction of C_t -N, and therefore an increase in the spin-marker RR frequencies. (We were unfortunately unable to obtain RR spectra for perchlorate adducts because of their instability upon laser excitation.) Lowering the pH of Fe¹¹¹ cytochrome c' from 10.3 to 6.9 decreases the magnetic moment appreciably⁴⁴ (reflecting a decreased S = $\frac{5}{2}$ contribution to the spin state⁴⁵) and also shifts the spin-marker RR frequencies⁴¹ up from typical five-coordinate values toward the values observed for low-spin Fe¹¹¹ (Table 11). We interpret the change in cytochrome c' structure between pH 10.3 and 6.9 as a weakening of the bond to the fifth ligand, accompanied by a strengthening of the Fe-pyrrole bonds, and a decrease in C_1 -N.

Core Expansion and Pyrrole Tilting. A. Linear Dependence on C_t -N. The present results establish that the correlation between the spin-marker frequencies and C_t -N discovered by Spaulding et al.¹¹ for band IV applies in a general way to iron porphyrin complexes. Huong and Pommier⁴⁷ have plotted the band IV and V frequencies against C_t -N for a large number of metalloporphyrins, and have found well-behaved linear relationships:

$$\nu = K(A - d) \,\mathrm{cm}^{-1}$$
 (1)

where d is the C_t-N distance and the parameters K (cm⁻¹/Å) and A (Å) have the values 555.6 and 4.86 for band IV and 423.7 and 5.87 for band V. Scholler and Hoffman²⁴ found a similar straight line for band IV. These lines are drawn in Figure 3 along with a third line which we have constructed to fit the available data for band 11, with K = 375.5 and A = 6.01.

The filled circles in Figure 3 are data points for a variety of

porphyrins with differing C_1 -N. Most of them fall satisfactorily on the lines. For $(1m)_2Fe^{11}MP$, $(2Melm)Fe^{11}MP$, and $(Me_2SO)_2Fe^{111}PP^+$, the frequencies fall on the band 1V line, but significantly $(15-20 \text{ cm}^{-1})$ below the lines for bands II and V. These latter bands have previously been noted as being sensitive to oxidation as well as spin state in iron porphyrins.^{2,8}

Bands 11, 1V, and V show substantial (4-11 cm⁻¹) shifts upon methine deuteration, whereas other bands in the 1350-1650-cm⁻¹ region do not.^{8,11} The three normal modes are thus assignable primarily to the stretching of the methine bridge bonds. The inverse correlation with C_{1} -N is plausibly linked to the decrease in the associated force constant when the methine bridges are stretched and bent to accommodate core expansion. The lowered band II and V frequencies observed for low-spin Fe¹¹ reflect its propensity for back-donation of d_{π} electrons to π^* orbitals of the porphyrin ring; the frequencies are raised when the axial imidazole ligands are replaced by good π acceptors, which compete with the porphyrin for the iron d_{π} electrons. Because the lowest porphyrin π^* orbital has nodes between the nitrogen atoms and the methine carbon atoms,48 back-donation is not expected to influence directly the methine bridge force constants, but rather those internal to the pyrrole ring. The internal coordinates are mixed in the normal modes, ^{10,49-51} however, and the influence of back-donation on the band 11 and V frequencies may be a reflection of the extent of mixing. The lack of sensitivity of band IV may be related to the fact that the mode is of A_{2g} symmetry, and stretching of the pyrrole C_{β} - C_{β} pyrrole bonds cannot contribute to the A_{2g} block. We are unable to explain the anomalously low band 11 and V frequencies for $[(Me_2SO)_2Fe^{11}PP]^+$ or for (2Melm)Fe¹¹MP, however.

B. Doming and Ruffling. The open circles in Figure 3 are data points for nonplanar porphyrins. They fall systematically below the lines, and we attribute the deviations to lowering of the methine bridge bond-stretching force constants resulting from reduction in π overlap when the pyrrole rings are tilted out of the mean plane. There are two kinds of pyrrole tilt which are observed in metalloporphyrin structures and have been discussed by Hoard.9 They are shown diagrammatically in Figure 4. In ruffled porphyrins the pyrrole rings swivel about the metal-nitrogen bonds, reducing the symmetry to D_{2d} . This leaves a square-planar MN₄ array, but permits a shortening of the M-N bond.⁹ The ruffled structure is observed for metals whose M-N bonds are shorter than Ct-N for an unconstrained planar porphyrin, estimated by Hoard⁹ to be 2.01 Å. An interesting example is provided by NiOEP, which can be crystallized in two forms. The triclinic modification⁵² contains a planar porphyrin, with Ni-N = 1.958 Å, while the orthorhombic form⁵³ contains a ruffled porphyrin with Ni-N =1.929 Å. There is evidently a close balance between loss of π conjugation in the ruffled structure and loss of Ni-N bond strength in the planar structure, which can be shifted by crystal packing forces. Spaulding et al.¹¹ observed that the two forms gave distinctly different RR spectra. While band IV was at the frequency predicted by the Ct-N correlation for the planar form, it decreased appreciably in the ruffled form, contrary to the Ct-N correlation. Ruffled structures are also found for four-coordinate Fe¹¹TPP¹³ and for [(1m)₂Fe¹¹¹OEP]Cl,⁵⁴ in both of which the Fe-N distance is significantly shorter than 2.01 Å.

A second type of pyrrole tilting is found in domed porphyrins, which are usually five coordinate.⁹ The metal atom lies above the porphyrin plane and this motion is followed by the pyrrole rings, which swivel about the methine bridge bonds causing the nitrogen plane, P_N , to lie above the mean plane of the porphyrin, P_C . The symmetry reduces to C_{4v} .

C. A Quantitative Model. As demonstrated by the crosses in Figure 3, the frequencies of the nonplanar porphyrins are



Figure 4. Perspective drawings for (a) planar, (b) domed, and (c) ruffled porphyrin skeletons from $\sim 10^{\circ}$ above. θ is the angle by which the pyrroles tilt out of the mean porphyrin plane.

accurately calculated by the equation

$$\nu = (K^2 (A - d)^2 \cos^2 \phi + B^2 \sin^2 \phi)^{1/2}$$
(2)

where the parameter *B* (cm⁻¹) is 1332, 1347, and 1438 for bands 11, 1V, and V. ϕ is the methine bridge dihedral angle between the N-C_{α}-C_m and C_{α}-C_m-C_{α} planes. Its calculation from the crystallographic data is described in Appendix A. For planar porphyrins $\phi = 0$, and eq 2 reduces to eq 1.

The derivation of eq 2 is given in Appendix B. It is based on a model of uncoupled normal modes associated solely with $C_{\alpha}-C_m$ bond stretching, with force constants made up of σ and π contributions, F_{σ} and F_{π} . The influence of core expansion and nonplanarity is carried by F_{π} . From the empirically determined parameters, K. A, and B, for bands 11, IV, and V, respectively, this model gives $F_{\sigma} = 7.7, 5.4$, and 6.2 mdyn/Å and $F_{\pi} = 2.1, 1.3$, and 1.8 mdyn/Å for an unconstrained porphyrin with C_t -N = 2.01 Å.⁹ Neglect of coupling with other normal modes could well account for the differences among the three estimates.

Taken together, however, these values are not far from what one might expect for methine bridges of bond order ~ 1.5 , considering that bond stretching force constants are roughly 5 and 10 mdyn/Å for single and double carbon-carbon bonds.⁵⁸ They suggest that the present model of combined core expansion and pyrrole tilting effects is not completely unreasonable, and could serve as a starting point for a more sophisticated treatment.

While the good account of the data given by eq 2 is satisfying, it also establishes that there is no unique way of using the RR frequencies to determine either core size or pyrrole tilt in porphyrins of unknown structure. Each set of frequencies is consistent with a range of C_t -N and θ values. In Figure 5 the loci of these values are shown for different values of the band IV frequency. The curves are quite flat for moderate values of θ , demonstrating that core size is the main determinant of the vibrational frequencies. If severe tilting were to occur, however, it would produce the same lowering of the frequencies as would modest core expansion.

Conclusions

At present there is no RR evidence to suggest significant protein-induced distortion of the porphyrin core in the heme proteins so far studied. In all cases their derivatives can be modeled with reasonable accuracy by simple porphyrin complexes. It is clear that the number of axial ligands and their properties are the main determinants of the porphyrin vibrational frequencies. Their influence on the spin-marker frequencies is related to porphyrin core expansion through steric interactions (although electronic effects are also evident for low-spin Fe¹¹). In addition porphyrin ruffling or doming



Figure 5. Loci of C_t -N distance and pyrrole tilt angles, allowed by eq 2 and the geometric relations in the Appendix for the band IV frequencies of the indicated complexes. Symbols and structural references as in Figure 3.

have subsidiary effects on the frequencies. Since domed porphyrins (five-coordinate Fe^{11} and Fe^{11}) show essentially the same frequencies whether inside a heme protein or not, there is not reason to think that proteins exert a marked effect on doming, as had been suggested for metHb and metMb.⁸

It should be emphasized, however, that the present treatment is highly simplified. There are certainly more distortion coordinates that might play a role than a single core size and a single methine dihedral angle. We do not understand at present why both $(2Melm)Fe^{11}MP$ and $[(Me_2SO)_2Fe^{11}PP]^+$ show the expected frequency for band IV, but low values for bands II and V, or why fluorometHb and Mb shows a decrease in band IV, but not bands II and V, relative to aquometHb and Mb or to $[(Me_2SO)_2Fe^{11}PP]^+$. This behavior suggests additional methine bridge influences which distinguish among bands II, IV, and V.

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Appendix A. Methine Dihedral Angles

Pyrrole tilting produces a rotation about the $C_{\alpha}C_{m}$ bonds, which is equal and opposite for the two bonds at each methine bridge. The difference between domed and ruffled porphyrins is that pairs of rotations have the same phase sequence for each successive methine bridge in the domed structure, but alternate in phase in the ruffled structure. The connection between the $C_{\alpha}C_{m}$ rotation (dihedral) angles, ϕ , and the pyrrole tilt angle, θ , is given by the following trigonometric relations:

doming:
$$\phi_d = \sin^{-1} (\tan \psi_d \tan (\gamma/2)) + \sin^{-1} (\tan \psi_d \tan \tau)$$

ruffling:
$$\phi_r = \sin^{-1} (\tan \psi_r \tan (\gamma/2))$$

- $\sin^{-1} (\tan \psi_r \cot \tau)$

The angles γ and τ are defined by

$$\gamma = \angle (C_{\alpha}C_{m}C_{\alpha})$$
$$\tau = \angle (C_{m}C_{\alpha}N) - \frac{1}{2}\angle (C_{\alpha}NC_{\alpha})$$

and θ is the angle of tilt of the pyrrole rings with respect to the mean porphyrin plane. In the present model, the pyrrole rings are assumed to be coplanar with the adjacent C_m atoms, and the methine bridge angles are given by

doming:
$$\cos \gamma = |1.0 - (\sin \tau + \cos \tau \cos \theta)^2|$$

ruffling: $\cos \gamma = |1.0 - (\cos \tau + \sin \tau \cos \theta)^2|$

Actual crystal structures do not necessarily conform strictly to the idealized C_{4v} or D_{2d} symmetries assumed for domed and ruffled porphyrins. An average tilt angle, θ_{av} , can be calculated, from the reported atomic displacements from the mean plane.⁵² For domed porphyrins θ_{av} is calculable from the reported P_N-P_C distance.⁵⁵ For porphyrin structures we have examined the angle τ is essentially constant, 72.5 ± 0.5°. The parameters ψ_d and ψ_r are given by

$$\sin\psi_{\rm d}=\cos\tau\,\sin\theta$$

 $\sin\psi_{\rm r}=\sin\tau\,\sin\theta$

Appendix B. Interpretation of Equation 2

The vibrational secular equation is given by $|\mathbf{GF} - \lambda| = 0$

$$-\lambda|=0 \tag{3}$$

where G and F are the kinetic and potential energy matrices 57 and

$$\lambda = 4\pi^2 c^2 \nu^2 \tag{4}$$

(c is the velocity of light and ν is expressed in cm⁻¹). For the present model we assume that the normal modes are not coupled and that the solution

$$\lambda = GF \tag{5}$$

applies to each of them, G and F being the appropriate (diagonal) matrix elements. The normal mode force constant, F, may be arbitrarily divided into contributors from σ and π bonding:

$$F = F_{\sigma} + F_{\pi} \tag{6}$$

It seems reasonable to attribute tilting effects to F_{π} which should vary as some function of $\cos \phi$. We select

$$F_{\pi} = F_{\pi}^{0} \cos^2 \phi \tag{7}$$

because it gives reasonable values of F_{π}^{0} ; the fact that there are actually two adjacent dihedral angles at each methine bridge, which are equal but have opposite senses, lends some plausibility to the use of $\cos^2 \phi$ rather than $\cos \phi$. We also choose to place the influence of core expansion on F_{π}^{0} . The reason for this is that the experimental data given in Figure 3 indicate that the effect of ring tilting decreases with increasing core size; this interaction can be taken into account by allowing the frequency decrease on core expansion to reflect an associated decrease in F_{π} , keeping F_{σ} constant.

Combining 1, 4, 5, and 6 we have for planar porphyrins that

$$K^{2}(A-d)^{2} = G'(F_{\sigma} + F_{\pi}^{0})$$
(8)

where $G' = G/4\pi^2 c^2$, and that

$$F_{\pi}^{0} = K^{2}(A - d)^{2}/G' - F_{\sigma}$$
(9)

With the above assumptions the relationship for tilted porphyrins should be

$$\nu = (G'(F_{\sigma} + F_{\pi}^{0} \cos^{2} \phi))^{1/2}$$
(10)

which, on substitution of eq 9, becomes

$$\nu = (K^2 (A - d)^2 \cos^2 \phi + G' F_\sigma \sin^2 \phi)^{1/2}$$
(11)

Equation 11 gives eq 2 with the identification $B = \sqrt{G'F_{\sigma}}$.

If the normal modes consisted solely of C_{α} - C_m stretching, then the G matrix elements would simply be $(1 + 2 \sin^2 \gamma)/m_c$ $(m_{\rm c}, \text{the carbon mass, is 12})$ for both band IV and V (the A_{2g} and B_{1g} symmetries both involve out-of-phase stretching of adjacent C_{α} - C_{m} bonds), and $(1 + 2\cos^{2}\gamma)/m_{c}$ for band 11 (A_{1g}, in phase), where γ is the $C_{\alpha}C_{m}C_{\alpha}$ angle, 124°. These values of G were combined with the empirically determined values of B to calculate the F_{σ} values given in the text. Equation 9 gives F_{π}^{0} , which depends on d; the estimate of 2.01 Å for an unconstrained porphyrin is due to Hoard.9

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